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► To cite this version:

Shyam Sundar Ghoshal, Billel Guelmame, Animesh Jana, Stéphane Junca. Optimal regularity for all time for entropy solutions of conservation laws in BV^s . *Nonlinear Differential Equations and Applications*, 2020, 27 (5), pp.46. 10.1007/s00030-020-00649-5 . hal-02495036v2

HAL Id: hal-02495036

<https://hal.science/hal-02495036v2>

Submitted on 13 Aug 2020

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Optimal regularity for all time for entropy solutions of conservation laws in BV^s

Shyam Sundar Ghoshal*, Billel Guelmame[†],
Animesh Jana[‡], Stéphane Junca[§]

Abstract

This paper deals with the optimal regularity for entropy solutions of conservation laws. For this purpose, we use two key ingredients: (a) fine structure of entropy solutions and (b) fractional BV spaces. We show that optimality of the regularizing effect for the initial value problem from L^∞ to fractional Sobolev space and fractional BV spaces is valid for all time. Previously, such optimality was proven only for a finite time, before the nonlinear interaction of waves. Here for some well-chosen examples, the sharp regularity is obtained after the interaction of waves. Moreover, we prove sharp smoothing in BV^s for a convex scalar conservation law with a linear source term. Next, we provide an upper bound of the maximal smoothing effect for nonlinear scalar multi-dimensional conservation laws and some hyperbolic systems in one or multi-dimension.

Key words: Conservation laws, entropy solutions, shocks, smoothing effect, fractional BV spaces BV^s .

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1 Introduction

For nonlinear conservation laws, it is known since Lax-Oleřnik [36, 41] that the entropy solution can have a better regularity than the initial data for Burgers type fluxes. Such smoothing effect has been obtained in fractional Sobolev spaces [37] and recently in fractional BV space [13] for more general fluxes. The optimality of such regularization is largely open in general. For scalar 1-D conservation laws, there are some optimal results proven up to finite time [18, 25, 38]. The aim of this article is to obtain the same optimality for all time.

We start with the one-dimensional scalar conservation laws which reads as follow:

$$\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} = 0 \quad \text{for } x \in \mathbb{R}, t > 0, \quad (1)$$

$$u(x, 0) = u_0(x) \text{ for } x \in \mathbb{R}. \quad (2)$$

The classical well-posedness theory for the Cauchy problem (1)–(2) is available for L^∞ and BV initial data [35, 36, 41]. BV -regularizing effect on entropy solutions has been established in [36, 41] for uniformly convex fluxes. It is well known that if the flux function is not uniformly convex then in general, the entropy solution of (1) may not have a finite total variation, [4, 21]. It can be shown that in one dimension if f'' vanishes at some point then there exists a class of initial data such that f can not regularize the corresponding entropy solution up to BV for all time [28]. Hence, to understand the optimal regularity of the entropy solution of (1), one works with more general space like fractional Sobolev space $W^{s,p}$ and fractional BV spaces BV^s , $0 < s < 1$, $1 \leq p$.

The advantage of BV^s spaces is to recover the fractional Sobolev regularity $W^{\sigma,p}$ for all $\sigma < s$, $1 \leq p < s^{-1}$ and to get the BV like trace properties of entropy solutions [23, 42, 43]. In one dimension, existence of the entropy solutions of (1) in BV^s , with BV^s data has been done in [13] and with L^∞ data in the same spaces in [13, 38, 39]. For non-convex fluxes a Lagrangian framework is used [8, 9]. For the scalar 1-D case, the BV^s smoothing effect corresponds to the optimal smoothing effect conjectured by Lions, Perthame and Tadmor in Sobolev spaces with the same fractional derivative s [37]. In multi-dimension, for a $C^{2,\gamma}$ flux, it has been shown [28] that entropy solutions do not need to have fractional derivative $s + \varepsilon$ for $\varepsilon > 0$. For multi-dimensional scalar conservation laws, regularizing effect in fractional Sobolev space was first studied in [37]. We refer [44] for the best known result in this direction and [31] for further improvement with some extra assumptions, see also [26] for such results with a source term. The proof of optimality of the exponent $s > 0$ is limited to some one-dimensional scalar examples [18, 25] and before the nonlinear interaction of waves. It has been extended for the scalar multi-dimensional case in [20, 32] but not for all time. Recently, in [28] it has been shown that in multi-dimension for any C^2 flux f there exists initial data u_0 such that the corresponding entropy solution is not in BV for all time $t > 0$.

The present article resolves the following:

- In one dimension, the optimal smoothing effect in fractional BV spaces is known for the equation (1) in bounded strip of time $(0, T)$ for $T > 0$, before the interactions of waves [18]. So it is natural to ask the following question:

Does there exists an entropy solution to (1) with compact support such that it does not belong to $BV^{s+\varepsilon}$ for all $t > 0$, $\varepsilon > 0$? (Q)

where s depends on the non-linearity of flux function. We first obtain an entropy solution to (1) for flux $f(u) = (1+p)^{-1}|u|^{1+p}$ such that $TV^{s+\varepsilon}(u(\cdot, t)) = \infty$, for all $\varepsilon > 0$, for all $t > 0$,

whereas $TV^s(u(\cdot, t)) < \infty$ with $s = p^{-1}$. Later we generalize this result for a larger class of fluxes.

- We extend the above result to higher dimension under some smooth regularity assumption on the flux in section 4.
- We are also able to answer the question (Q) for entropy solutions to balance laws which read as follow where $\alpha \in L^\infty((0, +\infty), \mathbb{R})$:

$$\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} = \alpha(t)u \quad \text{for } x \in \mathbb{R}, t > 0, \quad (3)$$

$$u(x, 0) = u_0(x) \quad \text{for } x \in \mathbb{R}. \quad (4)$$

- Smoothing effect for balance laws of type (3) in fractional BV space is not known. Based on a recent Lax-Oleinik type formula [5] we prove the BV^s regularizing effect for entropy solutions to (3) with a convex flux satisfying the p -degeneracy power law condition [13],

$$\frac{|f'(u) - f'(v)|}{|u - v|^p} \geq c_0 > 0 \text{ for } u \neq v \in [-M, M]. \quad (5)$$

We recall that if $f \in C^2([-M, M])$ then $p \geq 1$ [13]. The exact power-law degeneracy is given by the infimum of p satisfying (5). When f is smooth, the infimum is a minimum [13].

- Finding an upper bound of the maximal regularity, even in a given strip $(\mathbb{R} \times [0, T])$, was unknown for some triangular systems of conservation laws e.g., a pressure swing adsorption system. We answer this question for a class of 1-D triangular systems and the multi-D Keyfitz-Kranzer system in section 5 and 6 respectively.

To provide an answer to the question (Q) we recall some of the previously constructed examples [2, 3, 4]. In Section 3, Theorem 3.1 provides the direct answer to (Q) for a power-law type flux function $f(u) = (1 + p)^{-1}|u|^{1+p}$. We have discussed before that convex flux function with p -degeneracy (i.e., satisfying (5)) gives a regularizing effect in BV^s with $s = 1/p$ [13]. We construct an entropy solution u to (1) such that $TV^{s+\varepsilon}(u(\cdot, t)) = \infty$ for all $t > 0$ and $\varepsilon > 0$ with $s = 1/p$. Following the constructions in [3, 4] we build the entropy solution u , consisting infinitely many shock profiles in a compact interval. These shock profiles are named *Asymptotically Single Shock Packet (ASSP)* in [2]. Loosely speaking an ASSP is a solution with a special structure between two parallel lines in the half plane $\mathbb{R}_x \times \mathbb{R}_t^+$ such that in large time only one shock curve appears between them. ASSP plays a role to describe structure and large time behaviour of the entropy solution to strictly convex flux [2]. For the more complex structures of solutions for non-convex fluxes we refer interested reader to [8, 9]. The construction is done in Section 3. The building block of such solutions has a support in half strip $[a, b] \times \mathbb{R}_+$ for some $a < b$ and having an oscillation of amount δ_n up to time t_n . Then we club all of these building-blocks to get a solution with the same regularity for all time. Similar type of constructions for a slightly different aspect have been used in [1, 27] to build non BV solutions of scalar conservation laws with discontinuous fluxes. A larger class of non-uniformly convex fluxes has been considered in [4] to build non- BV solutions for all time. Such a flux f satisfies,

$$0 < f'(a) - f'(b) \leq C(a - b)^q \text{ for all } b < \theta_f < a \quad (6)$$

where $q > 1$ and $C > 0$, which implies that $f''(\theta_f) = 0$. Condition (6) is about a minimal degeneracy condition of f' near the point where f'' vanishes. We answer the question (Q) for this general class of convex functions satisfying (6).

The notion of *ASSP* has been generalized recently in [5] for balance laws of type (3). Based on a Lax-Oleřnik type formula [5] for entropy solutions to (3), we are able to answer analogous version of (Q) for such a balance law with linear source term. In Section 3 we provide a construction to show the optimality of the regularizing effect in balance law set up for power-law type flux functions. Like the case for $\alpha \equiv 0$, the constructed solution for balance law is a juxtaposition of infinitely many *ASSP*. Naturally, for such balance laws, the boundaries of *ASSP* are curves instead straight lines. Moreover, when $\alpha \equiv 0$, the case of conservation laws is recovered. We choose to answer to the question (Q) in this slightly more general setting.

In the remainder of the paper, Sections 4, 5 and 6, the results obtained for the one dimensional scalar case are used to bound the maximal smoothing effect for solutions of three different problems, namely, scalar multidimensional equations, a class of triangular systems and a multidimensional system. For the multidimensional case, planar waves are used as in [20, 28, 32]. For a class of triangular systems involving a transport equation, the main problem is to keep the linear component bounded and not being a δ -shock while the nonlinear component belongs to the critical BV space. For multi-dimensional Keyfitz-Kranzer system [34], it has been shown, [24] that small *TV* bound of initial data is not enough to get immediately a *BV* renormalized solution of the Keyfitz-Kranzer system. In this article we implement his construction to get a similar blow up in all BV^s spaces, $s > 0$.

2 Fractional BV spaces, BV^s , $0 < s \leq 1$

In this section, the definition of generalized $BV(\mathbb{R})$ spaces are recalled [40]. Then the multi-D case is stated.

Definition 2.1 ($BV^s(\mathbb{R}, \mathbb{R})$). *Let $p = 1/s$, the TV^s variation also called the total p -variation of any real function v is:*

$$TV^s v = \sup_{\{x_i\} \in \mathcal{P}} \sum_{i=2}^n |v(x_i) - v(x_{i-1})|^p \quad (7)$$

where $\mathcal{P} = \{\{x_1, \dots, x_n\}, x_1 < \dots < x_n, 2 \leq n \in \mathbb{N}\}$ is the set of subdivisions of \mathbb{R} .

The space $BV^s(\mathbb{R}, \mathbb{R})$ is the subset of real functions such that,

$$BV^s(\mathbb{R}) = \{v, TV^s(v) < \infty\}. \quad (8)$$

Notice that $BV^1 = BV$ and $BV^s \subset L^\infty$ for all $0 < s \leq 1$. By convention, we set $BV^0 = L^\infty$. A similar definition can be used to defined $BV^s(I, \mathbb{R})$ where $I \subset \mathbb{R}$, only considering the subdivisions of I . The fractional Sobolev space $W^{s,p}$ can be defined as follows:

Definition 2.2. *Let $\Omega \subset \mathbb{R}^N$ be open. Let $s \in (0, 1)$ and $p \in [1, \infty)$. By $W^{s,p}(\Omega)$ we denote the set of all $u \in L^p(\Omega)$ such that*

$$\frac{|u(x) - u(y)|}{|x - y|^{s + \frac{N}{p}}} \in L^p(\Omega \times \Omega). \quad (9)$$

It is worth mentioning that BV^s does not coincide with fractional Sobolev space, $W^{s,p}$ but it is closely related to $W^{s,p}$ with the critical exponent p for the traces theory, that is, $sp = 1$. More precisely, for all $\varepsilon > 0$, $BV_{loc}^s \subset W_{loc}^{s-\varepsilon, 1/s} \subset W_{loc}^{s-\varepsilon, 1}$ [13]. All the examples valid for all times in this article present shocks, so are discontinuous and therefore never belong to $W^{s+\varepsilon, 1/s}$, $\forall \varepsilon > 0$.

Thus, a non BV^s regularity corresponds to a non Sobolev regularity with the same exponent up to any positive ε . The optimality can also be studied in BV^s and corresponds to the similar Sobolev optimaty. Notice also that the estimates in fractional BV spaces can be simpler than in fractional Sobolev spaces as in [17]. It is the reason why only result in BV^s are given in this paper.

Furthermore, BV^s regularity guarantees left and right traces like BV functions. That is why BV^s spaces seem more well fitted to study the regularity of the solutions of conservation laws than the corresponding Sobolev spaces.

To extend the definition of BV^s for the multi-D case, a theorem characterizing BV^s in 1-D is used. This theorem characterizes the space BV^s with the Holder space Lip^s and the BV space. It is due to Michel Bruneau [16].

Theorem 2.1 (Bruneau, 1974). *For any $u \in BV^s$ there exists the following factorization by a s -Holder function and a BV function,*

$$u \in BV^s \Leftrightarrow \exists L \in \text{Lip}^s(\mathbb{R}, \mathbb{R}), \exists v \in BV(\mathbb{R}) \text{ s.t. } u = L \circ v.$$

That means that

$$BV^s(\mathbb{R}, \mathbb{R}) = \text{Lip}^s(\mathbb{R}, \mathbb{R}) \circ BV(\mathbb{R}, \mathbb{R}).$$

In order to define $BV^s(\mathbb{R}^m)$, we recall the definition of $BV(\mathbb{R}^m)$ for $m \geq 1$

Definition 2.3 ($BV(\mathbb{R}^m)$). *A function u belongs to $BV(\mathbb{R}^m)$ if there exists a Radon measure μ such that*

$$\int_{\mathbb{R}} u(x) \operatorname{div} \phi(x) \, dx = - \langle \mu, \phi \rangle \quad \forall \phi \in \mathcal{C}_c^1(\mathbb{R}^m).$$

Now, the following natural definition of $BV^s(\mathbb{R}^m)$ is proposed for $m \geq 1$.

Definition 2.4 ($BV^s(\mathbb{R}^m)$). *A function u belongs to $BV^s(\mathbb{R}^m)$ if there exists the following factorization by an s -Holder function $L \in \text{Lip}^s(\mathbb{R}, \mathbb{R})$ and a $BV(\mathbb{R}^m, \mathbb{R})$ function v such that*

$$u = L \circ v.$$

That means that

$$BV^s(\mathbb{R}^m, \mathbb{R}) = \text{Lip}^s(\mathbb{R}, \mathbb{R}) \circ BV(\mathbb{R}^m, \mathbb{R}). \quad (10)$$

This definition can be extended to $BV_{loc}^s(\mathbb{R}^m)$ by:

$$BV_{loc}^s(\mathbb{R}^m, \mathbb{R}) = \text{Lip}^s(\mathbb{R}, \mathbb{R}) \circ BV_{loc}(\mathbb{R}^m, \mathbb{R}). \quad (11)$$

Notice that the Holder function has to be globally on \mathbb{R} an Holder function since $BV(\mathbb{R}^m)$ is not a subset of L^∞ for $m > 1$.

This definition is valid for $m = 1$ thanks to Bruneau's Theorem 2.1. Moreover, a $BV_{loc}^s(\mathbb{R})$ 1-D function can be also considered as a $BV_{loc}^s(\mathbb{R}^m)$ multi-D function by the following lemma.

Lemma 2.1. *Let $\xi \in S^{m-1}$ and $U(X) = u(\xi \cdot X)$, $U \in BV_{loc}^s(\mathbb{R}^m)$ if and only if $u \in BV_{loc}^s(\mathbb{R})$*

Proof. From the Bruneau's Theorem 2.1, slightly extended on bounded set, $u(x) = L(v(x))$ where $v \in BV_{loc}(\mathbb{R})$. Let $V(X)$ be $v(\xi \cdot X)$. V belongs to $BV_{loc}(\mathbb{R}^m)$ [6, 29]. Thus $U(X) = L(V(X))$ belongs to $BV_{loc}^s(\mathbb{R}^m, \mathbb{R})$. The converse is also clear. \square

3 Sharp regularity for scalar 1D entropy solutions

In this section, we will build some examples to show the optimality of smoothing effect in BV^s for all time. This regularity has been obtained in [13, 30, 17, 38, 39]. The optimality for all time is new. For that purpose, we consider the flux $f(u) = |u|^{p+1}/(p+1)$ so $f'(u) = u|u|^{p-1}$. It is shown that for $u_0 \in L^\infty$, the solution becomes instantly in BV_{loc}^s , with $s = p^{-1}$. Theorem 3.1 stated below shows that the regularizing in BV^s space is optimal for all time since there exist entropy solutions u such that $u(\cdot, t) \notin BV^{s+\varepsilon}$ for all $\varepsilon > 0$ and for all $t > 0$. The construction of this example is similar to the one done in [2] to show infinitely many shock profile occurrence in compact interval. Similar construction has been also used in [4] to show the existence of an entropy solution which does not belong to BV for all time. Here we use it to show the existence of an entropy solution which is exactly in BV^s with $s = p^{-1}$ for all time $t > 0$ with no more regularity.

Theorem 3.1. *There exists compactly supported initial data $u_0 \in L^\infty(\mathbb{R})$ such that the corresponding entropy solution $u(\cdot, t) \in L^\infty(\mathbb{R} \times [0, \infty))$ of the scalar conservation law (1) with the flux $f(u) = |u|^{p+1}/(p+1)$, $p \geq 1$ satisfies for all $t > 0$, for all $\varepsilon > 0$ with $s = p^{-1}$,*

$$TV^s u(\cdot, t) < +\infty = TV^{s+\varepsilon} u(\cdot, t).$$

Theorem 3.1 can be seen as a particular case (that is, $\alpha \equiv 0$) of the following result stated in context of balance laws.

Theorem 3.2. *There exists an initial data $u_0 \in L^\infty(\mathbb{R})$ such that the corresponding entropy solution to balance law (3) with flux $f(u) = (1+p)^{-1}|u|^{1+p}$ for $p > 1$ and $\alpha \in L^\infty(0, \infty)$ satisfies the following with $s = p^{-1}$*

$$TV^s(u(\cdot, t)) < \infty = TV^{s+\varepsilon}(u(\cdot, t)) \text{ for all } t > 0 \text{ and } \varepsilon > 0. \quad (12)$$

Theorem 3.2 also states about the regularizing of entropy solution corresponding to a particular initial data u_0 and flux $f(u) = (p+1)^{-1}|u|^{p+1}$. Next we will show that it is not restricted to a special choice of data and flux. If a flux satisfies a p -degeneracy condition like (5) then regularizing is valid for any L^∞ initial data. More precisely, we have the following

Theorem 3.3. *Let $f \in C^1(\mathbb{R})$ be a convex flux satisfying a power-law condition (5) and super linear growth condition (15). Let $\alpha \in L^\infty(0, \infty)$. Let $u_0 \in L^\infty(\mathbb{R})$. Let u be the entropy solution of the initial value problem for the balance law (3), with the initial data u_0 (4), then*

$$u(\cdot, t) \in BV_{loc}^s(\mathbb{R}) \text{ for } s = \frac{1}{p} \text{ and } \forall t > 0. \quad (13)$$

As we have discussed before for $\alpha \equiv 0$ case, that for entropy solutions to (1), uniformly convex flux regularizes the solution in BV space [36, 41] and it fails once we drop the uniform convexity assumption on flux function [4, 21]. As a natural extension, one can ask for the regularizing effect for strictly convex fluxes and it has been shown in [13] that regularizing is valid in fractional BV space once the flux satisfying a p -degeneracy condition (5). For strictly convex Lipschitz flux, regularizing effect can be obtained in more general spaces like BV^Φ with a special choice of Φ , [30]. To prove $TV^{s+\varepsilon}(u(\cdot, t)) = \infty$ for all $t > 0$ we construct an entropy solution consisting ASSP's (see [2] for more detail on ASSP). The other part, that is, $u \in BV^s$ follows from [13] in the case of Theorem 3.1 that is, when u solves (1). But for balance law of type (3) no such result exists.

It can be proved in a similar fashion as it was done in [13] for conservation laws. We first give a brief sketch of the proof for $u \in BV^s$ where u is the entropy solution to balance law (3). In order to do that let us first recall some of the definitions and results from [5].

Definition 3.1. Let $\alpha \in L^\infty(0, \infty)$. Let β be primitive of α , that is,

$$\beta(t) = \int_0^t \alpha(\theta) d\theta. \quad (14)$$

Suppose that the flux f is having super-linear growth, that is,

$$\lim_{|v| \rightarrow \infty} \frac{f(v)}{|v|} = \infty. \quad (15)$$

We define $\Psi : \mathbb{R} \times \mathbb{R}_+ \rightarrow \mathbb{R}$ -function as follows:

$$x = \int_0^t f' \left(\Psi(x, t) e^{\beta(\theta)} \right) d\theta \text{ for each } x \in \mathbb{R}. \quad (16)$$

Note that Ψ in (16) is well-defined on $\mathbb{R} \times \mathbb{R}_+$ due to super-linear growth (15) of f (see [5]). For $\alpha \equiv 0$ and strictly convex C^1 flux f , the Ψ -function is nothing but $(f')^{-1}(x/t)$. As it is observed in [5], the map $x \mapsto \Psi(x, t)$ is increasing for strictly convex flux f .

Proposition 3.1. ([5]) Let $\alpha \in L^\infty(\mathbb{R}_+)$ and the flux $f \in C^1(\mathbb{R})$ satisfying (15). Let u be the entropy solution to (3) with initial data $u_0 \in L^\infty(\mathbb{R})$. Then u satisfies

$$u(x, t) = e^{\beta(t)} \Psi(x - y(x, t), t) \text{ for all } x \in \mathbb{R}, t > 0 \quad (17)$$

for some function y such that $x \mapsto y(x, t)$ is non-decreasing and β is defined as in (14). Moreover, for each $T > 0$ there exists a constant $C(T)$ such that

$$|x - y(x, t)| \leq C(T)t \text{ for all } t \in [0, T]. \quad (18)$$

Since Ψ is increasing in its first variable we have the following lemma.

Lemma 3.1. Let $f \in C^1$ be a convex flux satisfying the super linear growth condition (15) and power-law condition (5) with $p \geq 1$. Then for any $z_1, z_2 \in \mathbb{R}$ we have with $s = p^{-1}$,

$$|\Psi(z_1, t) - \Psi(z_2, t)| \leq \left(\frac{|z_1 - z_2|}{c_0 \gamma(t)} \right)^s, \quad (19)$$

with

$$\gamma(t) := \int_0^t e^{p\beta(\theta)} d\theta \text{ where } \beta \text{ is defined as in (14)}. \quad (20)$$

Proof. Fix two points $z_1, z_2 \in \mathbb{R}$. Without loss of generality, we can assume $z_1 > z_2$. Since Ψ is increasing in its first variable, we have

$$\Psi(z_1, t) e^{-\beta(t)} \geq \Psi(z_2, t) e^{-\beta(t)}. \quad (21)$$

Since f' satisfies p -degeneracy condition (5) and f' is continuous, that means that f' is monotone. Assume that f' is increasing, so the absolute values are skipped in (5) and,

$$f'(a) - f'(b) \geq c_0(a - b)^p \text{ for } a \geq b. \quad (22)$$

Therefore, by the definition (16) of Ψ we also have for $z_1 > z_2$,

$$|z_1 - z_2| = z_1 - z_2 = \int_0^t f'(\Psi(z_1, t)e^{\beta(\theta)}) d\theta - \int_0^t f'(\Psi(z_2, t)e^{\beta(\theta)}) d\theta \quad (23)$$

$$\geq \int_0^t c_0(\Psi(z_1, t) - \Psi(z_2, t))^p e^{p\beta(\theta)} d\theta \quad (24)$$

$$= c_0|\Psi(z_1, t) - \Psi(z_2, t)|^p \int_0^t e^{p\beta(\theta)} d\theta. \quad (25)$$

This proves the inequality (19). \square

Now we are ready to prove the regularity result for entropy solution to balance laws (3).

Proof of Theorem 3.3. Fix a partition between $a = x_0 < x_1 < \dots < x_m = b$. By Proposition 3.1, we have

$$\sum_{k=1}^m |u(x_k, t) - u(x_{k-1}, t)|^p = e^{p\beta(t)} \sum_{k=1}^m |\Psi(x_k - y(x_k, t), t) - \Psi(x_{k-1} - y(x_{k-1}, t), t)|^p. \quad (26)$$

By virtue of Lemma 3.1 we have

$$\sum_{k=1}^m |u(x_k, t) - u(x_{k-1}, t)|^p \leq e^{p\beta(t)} (c_0 \gamma(t))^{-1} \sum_{k=1}^m |x_k - y(x_k, t) - x_{k-1} - y(x_{k-1}, t)|. \quad (27)$$

Since $x \mapsto y(x, t)$ is increasing for each fixed t , we have

$$\begin{aligned} \sum_{k=1}^m |u(x_k, t) - u(x_{k-1}, t)|^p &\leq e^{p\beta(t)} (c_0 \gamma(t))^{-1} [b - a + y(b, t) - y(a, t)] \\ &\leq e^{p\beta(t)} (c_0 \gamma(t))^{-1} [2(b - a) + 2C(T)t] \end{aligned} \quad (28)$$

for all $t \in (0, T)$. The last line follows from the inequality (18). This completes the proof of Theorem 3.3. \square

Our next aim is to establish the optimality of Theorem 3.3 for all time $t > 0$ and for that we restrict our discussion for power-law type fluxes, more precisely, $f(u) = (p + 1)^{-1}|u|^{p+1}$ for $p > 1$.

Proof of Theorem 3.2. To set the path for constructing an entropy solution to (3) which does not belong to $BV_{loc}^s(\mathbb{R})$ for $s > 1/p$, we first observe the structure of entropy solution to the following initial data.

$$u_0^n(x) = \begin{cases} 0 & \text{for } x < x_n - \Delta x_n, \\ \delta_n & \text{for } x_n - \Delta x_n < x < x_n, \\ -\delta_n & \text{for } x_n < x < x_n + \Delta x_n, \\ 0 & \text{for } x_n + \Delta x_n < x \end{cases} \quad (29)$$

for $\delta_n, \Delta x_n > 0$ and $x_n \in \mathbb{R}$. For $f(u) = (p+1)^{-1}|u|^{p+1}$, Ψ has the following form

$$\Psi(x, t) = x|x|^{-\frac{p-1}{p}}\gamma(t)^{-\frac{1}{p}} \quad (30)$$

where $\gamma(t)$ is defined as in (20). With the help of results from [2, 5] we have the following observations

1. Consider a Riemann problem w_C^0 defined as follows

$$w_C^0 = \begin{cases} w_- & \text{for } x < x_0, \\ w_+ & \text{for } x > x_0, \end{cases} \quad \text{where } w_- > w_+. \quad (31)$$

The entropy solution, w_C to (3) corresponding to Riemann data w_C^0 has the following form

$$w_C = \begin{cases} w_- e^{\beta(t)} & \text{for } x < x_0 + \lambda(t), \\ w_+ e^{\beta(t)} & \text{for } x > x_0 + \lambda(t), \end{cases} \quad (32)$$

for $t > 0$ where $\lambda(t)$ is defined as follows:

$$\lambda(t) := \frac{1}{w_+ - w_-} \int_0^t \left[f(w_+ e^{\beta(\theta)}) - f(w_- e^{\beta(\theta)}) \right] e^{-\beta(\theta)} d\theta. \quad (33)$$

2. Next we consider a special data w_L^0 defined as follows:

$$w_L^0(x) = \begin{cases} 0 & \text{for } x < x_L, \\ \delta_n & \text{for } x > x_L, \end{cases} \quad (34)$$

where $\delta_n > 0$. Then entropy solution to (3) with initial data w_L^0 will look like

$$w_L(x, t) = \begin{cases} 0 & \text{for } x < x_L, \\ \Psi(x - x_L, t) e^{\beta(t)} & \text{for } x_L \leq x \leq \zeta_L(t), \\ \delta_n e^{\beta(t)} & \text{for } x > \zeta_L(t), \end{cases} \quad (35)$$

for $t > 0$ where $\zeta_L(t)$ are determined by

$$\Psi(\zeta_L(t) - x_L, t) = \delta_n. \quad (36)$$

3. Now consider the following data

$$w_R^0(x) = \begin{cases} -\delta_n & \text{for } x < x_R, \\ 0 & \text{for } x > x_R, \end{cases} \quad (37)$$

where $\delta_n > 0$. Then entropy solution to (3) will look like

$$w_R(x, t) = \begin{cases} -\delta_n e^{\beta(t)} & \text{for } x < \zeta_R(t), \\ \Psi(x - x_R, t) e^{\beta(t)} & \text{for } \zeta_R(t) \leq x \leq x_R, \\ 0 & \text{for } x > x_R, \end{cases} \quad \text{for } t > 0, \quad (38)$$

where ζ_R is determined by

$$\Psi(\zeta_R(t) - x_R, t) = -\delta_n. \quad (39)$$

Let us set $\bar{x}_0 := x_n$, $x_L := x_n - \Delta x_n$ and $x_R := x_n + \Delta x_n$. Suppose the corresponding $\zeta_L(t)$ and $\zeta_R(t)$ intersect each other at (\tilde{x}_n, t_n) for the first time. From (35) and (38) we observe that $x_n - \Delta x_n = x_L \leq \tilde{x}_n \leq x_R = x_n + \Delta x_n$. By using (30), (36) and (39) we have

$$(\tilde{x}_n - x_n + \Delta x_n)^{\frac{1}{p}} \gamma(t_n)^{-\frac{1}{p}} = \delta_n = (x_n + \Delta x_n - \tilde{x}_n)^{\frac{1}{p}} \gamma(t_n)^{-\frac{1}{p}}. \quad (40)$$

Hence, we get $\tilde{x}_n = x_n$. Subsequently, we obtain $(\Delta x_n)^{\frac{1}{p}} \gamma(t_n)^{-\frac{1}{p}} = \delta_n$. Recall definition of $\gamma(t)$ as in (20). Therefore, t_n is determined as follows,

$$\int_0^{t_n} e^{p\beta(\theta)} d\theta = \frac{\Delta x_n}{\delta_n^p}. \quad (41)$$

Suppose B^* is the integration of e^β over \mathbb{R}_+ , that is,

$$B^* := \int_0^\infty e^{p\beta(\theta)} d\theta = \gamma(+\infty). \quad (42)$$

Note that $B^* \in (0, \infty]$. For $\alpha \equiv 0$ case, we have $B^* = \infty$. Next consider the case when $B^* < \infty$ and

$$\frac{\Delta x_n}{\delta_n^p} \geq B^*. \quad (43)$$

In this case, we have the following feature which does not arise for solutions of (1):

$$\zeta_L(t) < x_n < \zeta_R(t) \text{ for all } t \in (0, \infty). \quad (44)$$

Summarizing these observations we have

- O1.** If $\frac{\Delta x_n}{\delta_n^p} < B^*$ then (41) has a unique solution in $(0, \infty)$.
- O2.** If $\frac{\Delta x_n}{\delta_n^p} \geq B^*$ then (41) has no solution in $(0, \infty)$. In this case, we set $t_n = \infty$, that is to say that ζ_L and ζ_R never meet with each other.

If $t_n < \infty$ then note that for $t > t_n$ we have

$$\Psi(x_n - x_L, t) = -\Psi(x_n - x_R, t). \quad (45)$$

Therefore we have the following structure of entropy solution u , to (3) with initial data u_0 as in (29):

1. For $0 < t < t_n$ we have

$$u^n(x, t) = \begin{cases} 0 & \text{for } x < x_n - \Delta x_n, \\ \Psi(x - (x_n - \Delta x_n), t)e^{\beta(t)} & \text{for } x_n - \Delta x_n < x < \zeta_L(t), \\ \delta_n e^{\beta(t)} & \text{for } \zeta_L(t) < x < x_n, \\ -\delta_n e^{\beta(t)} & \text{for } x_n < x < \zeta_R(t), \\ \Psi(x - (x_n + \Delta x_n), t)e^{\beta(t)} & \text{for } \zeta_R(t) < x < x_n + \Delta x_n, \\ 0 & \text{for } x_n + \Delta x_n < x. \end{cases} \quad (46)$$

2. For $t > t_n$ we have

$$u^n(x, t) = \begin{cases} 0 & \text{for } x < x_n - \Delta x_n, \\ \Psi(x - (x_n - \Delta x_n), t)e^{\beta(t)} & \text{for } x_n - \Delta x_n < x < x_n, \\ \Psi(x - (x_n + \Delta x_n), t)e^{\beta(t)} & \text{for } x_n < x < x_n + \Delta x_n, \\ 0 & \text{for } x_n + \Delta x_n < x. \end{cases} \quad (47)$$

Note that $TV^{s+\varepsilon}(u^n(\cdot, t)) \geq (2\delta_n)^{\frac{1}{s+\varepsilon}} e^{\frac{\beta(t)}{s+\varepsilon}}$ for $t \in [0, t_n]$. From the above discussion we know that support of the entropy solution $u_n(\cdot, t)$ lies in $[x_n - \Delta x_n, x_n + \Delta x_n]$ for all time $t > 0$. We choose $\Delta x_n = (n \log^2(n+1))^{-1}$ and $\delta_n = (n \log^3(n+1))^{-\frac{1}{p}}$. Subsequently, we have

$$\frac{\Delta x_n}{\delta_n^p} = \log(n+1) \rightarrow \infty \text{ as } n \rightarrow \infty. \quad (48)$$

Since $\sum_{n=1}^{\infty} \Delta x_n < \infty$ we can choose a sequence x_n such that $x_n + \Delta x_n < x_{n+1} - \Delta x_{n+1} < x^* < \infty$ for all $n \geq 1$. Now we define an initial data u_0 as follows

$$u_0 = \sum_{n=1}^{\infty} u_0^n. \quad (49)$$

Note that by previous observation and choice of x_n , entropy solutions u_n has mutually disjoint support for all $t > 0$. Therefore, the entropy solution u of (3) corresponding to initial data u_0 can be written as

$$u(x, t) = \sum_{n=1}^{\infty} u_n(x, t) \text{ for all } x \in \mathbb{R}, t > 0. \quad (50)$$

Recall observations (O1.) and (O2.). Hence, for each fixed $t \in (0, \infty)$ there exists an n_0 such that $t < t_n$ for all $n \geq n_0$ due to (48). From definition (14) of $\beta(t)$ we have $\beta(t) \geq -t\|\alpha\|_{L^\infty(\mathbb{R}_+)}$ for all $t \geq 0$. Therefore, we have

$$TV^{s+\varepsilon}(u(\cdot, t)) \geq 2^{\frac{1}{s+\varepsilon}} e^{\frac{\beta(t)}{s+\varepsilon}} \sum_{n=n_0}^{\infty} \delta_n^{\frac{p}{1+p\varepsilon}} = 2^{\frac{1}{s+\varepsilon}} e^{\frac{-t}{s+\varepsilon}\|\alpha\|_{L^\infty(\mathbb{R}_+)}} \sum_{n=n_0}^{\infty} \frac{1}{(n \log^3(n+1))^{\frac{1}{1+p\varepsilon}}} = \infty. \quad (51)$$

Note that the $TV^s(u(\cdot, t)) < \infty$ for $s = 1/p$ and $t > 0$ due to Theorem 3.3. \square

Our next result upgrades Theorem 3.1 for more general class of functions satisfying the following hypothesis:

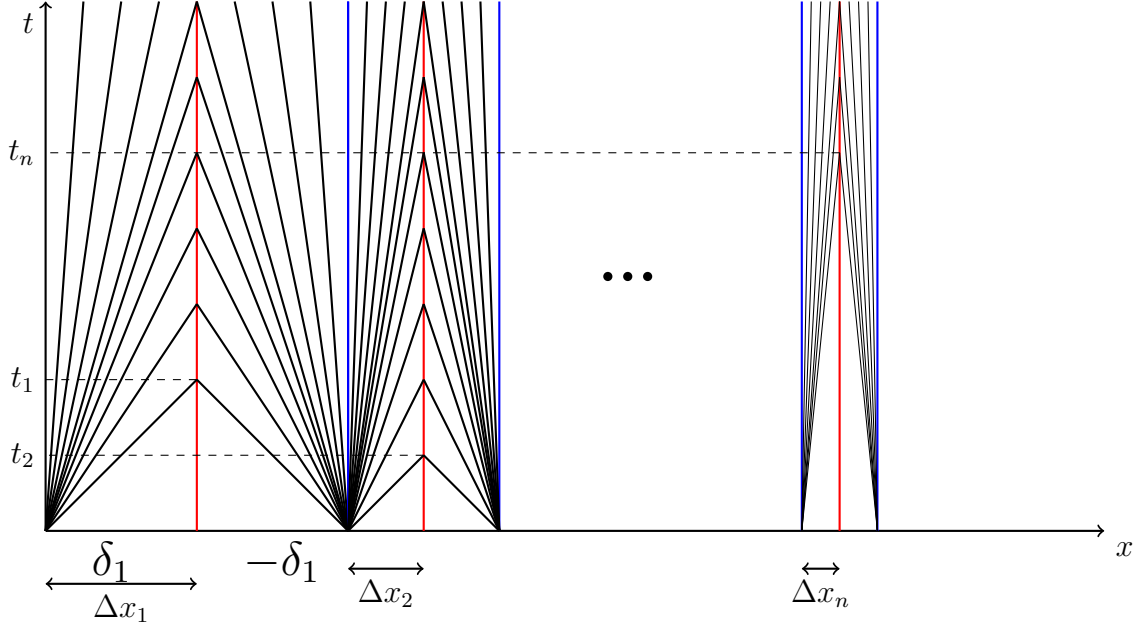


Figure 1: This picture illustrates the entropy solution constructed in Theorem 3.1 for $\alpha \equiv 0$ case. This construction and structure of entropy solution have been previously studied in [2, 3, 4].

(H-1) $f \in C^1(\mathbb{R})$ is a strictly convex function such that $f(0) = f'(0) = 0$.

(H-2) There exist $q > 1$, $r > 0$ and $C > 0$ such that

$$0 \leq f'(a) - f'(b) \leq C(a - b)^q \text{ for all } b \in (-r, 0) \text{ and } a \in (0, r). \quad (52)$$

Remark 3.1. We do not lose generality by assuming that $f(0) = f'(0) = 0$, due to the change of variables $x \mapsto x - f'(0)t$ and $\tilde{f}(u) \stackrel{\text{def}}{=} f(u) - f(0) - f'(0)u$.

The class of function satisfying (H-1) and (H-2) was previously considered in [4] to show non-BV propagation for all time $t > 0$. In this article, we will show the non- BV^s propagation for same class of function in the context of balance laws (3).

Theorem 3.4. Let $f \in C^1(\mathbb{R})$ satisfying (H-1) and (H-2) along with super linear growth condition (15). Let $\alpha \in L^\infty(0, \infty)$. Then there exists a compact support initial data $u_0 \in L^\infty(\mathbb{R})$ such that the corresponding entropy solution u of (3) satisfies the following:

$$u(\cdot, t) \notin BV_{loc}^s(\mathbb{R}) \text{ for all } s > 1/q, t > 0. \quad (53)$$

Remark 3.2. Note that Theorem 3.4 is optimal for the class of fluxes satisfying (52). It is easy to verify that $f(u) = (q+1)^{-1}|u|^{q+1}$ satisfies (52) and as we have seen in Theorem 3.3, $u(\cdot, t) \in BV^{1/q}$ for $t > 0$.

Observation: We want to make a remark that for convex f satisfying (H-1), we have $\Psi(0, t) = 0$. From the definition of Ψ we have

$$0 = \int_0^t f'(\Psi(0, t)e^{\beta(\theta)}) d\theta. \quad (54)$$

Since f is a C^1 strictly convex function, f' is increasing. Hence $af'(a) > 0$ for any $a \neq 0$ because $f'(0) = 0$. Suppose $\Psi(0, t) > 0$ then

$$\int_0^t f' \left(\Psi(0, t) e^{\beta(\theta)} \right) d\theta > 0. \quad (55)$$

Similarly, if $\Psi(0, t) < 0$, then we have

$$\int_0^t f' \left(\Psi(0, t) e^{\beta(\theta)} \right) d\theta < 0. \quad (56)$$

Note that both (55) and (56) contradict with (54). Therefore we have $\Psi(0, t) = 0$. Note that Ψ is increasing in its first variable due to strict convexity assumption on f . Subsequently, we get $x\Psi(x, t) > 0$ for any $x \neq 0$.

Before we give the main construction to prove Theorem 3.4 we first recall some results from [4] and find structure of the entropy solution to the following data

$$u_{A,B}^0(x) = \begin{cases} 0 & \text{if } x \notin [A, B], \\ a_{A,B} & \text{if } x \in [A, \tau], \\ b_{A,B} & \text{if } x \in (\tau, B]. \end{cases} \quad (57)$$

Next we make a choice for the pair $(a_{A,B}, b_{A,B})$ depending on A, B . For that purpose we define $G : \mathbb{R} \rightarrow \mathbb{R}$ as

$$G(a) = \int_0^{t_0} a f'(a e^{\beta(\theta)}) - f(a e^{\beta(\theta)}) e^{-\beta(\theta)} d\theta, \quad (58)$$

where $t_0 > 0$ is fixed. Now we claim that $z \mapsto G(z)$ is increasing for $z > 0$ and decreasing for $z < 0$. To see this consider $a > a_1 > 0$, then by Mean Value Theorem we have

$$\begin{aligned} & a e^{\beta(\theta)} f'(a e^{\beta(\theta)}) - a_1 e^{\beta(\theta)} f'(a_1 e^{\beta(\theta)}) - f(a e^{\beta(\theta)}) + f(a_1 e^{\beta(\theta)}) \\ &= a e^{\beta(\theta)} f'(a e^{\beta(\theta)}) - a_1 e^{\beta(\theta)} f'(a_1 e^{\beta(\theta)}) - f'(c^*) (a - a_1) e^{\beta(\theta)}, \end{aligned}$$

for some $c^* \in (a_1 e^{\beta(\theta)}, a e^{\beta(\theta)})$. Since f' is increasing, we get

$$\begin{aligned} & a e^{\beta(\theta)} f'(a e^{\beta(\theta)}) - a_1 e^{\beta(\theta)} f'(a_1 e^{\beta(\theta)}) - f(a e^{\beta(\theta)}) + f(a_1 e^{\beta(\theta)}) \\ & \geq a e^{\beta(\theta)} f'(a e^{\beta(\theta)}) - a_1 e^{\beta(\theta)} f'(a_1 e^{\beta(\theta)}) - f'(a e^{\beta(\theta)}) (a - a_1) e^{\beta(\theta)} \\ &= a_1 e^{\beta(\theta)} \left[f'(a e^{\beta(\theta)}) - f'(a_1 e^{\beta(\theta)}) \right] \\ & \geq 0. \end{aligned}$$

Hence, from (58) we obtain $a \mapsto G(a)$ is increasing for $a > 0$. By a similar argument, we get $b \mapsto G(b)$ is decreasing for $b < 0$. Therefore, we have that there exists $r_0 > 0$ such that for a given $a \in (0, r_0)$ there is a $b \in (-r, 0)$ satisfying $G(a) = G(b)$. Let us fix $a_0 \in (0, r_0)$ and $b_0 \in (-r, 0)$ such that $G(a_0) = G(b_0)$. Define $F_- : [b_0, 0] \rightarrow \mathbb{R}$ and $F_+ : [0, a_0] \rightarrow \mathbb{R}$ as follows

$$F_+(a) = \int_0^{t_0} f'(a e^{\beta(\theta)}) d\theta \text{ and } F_-(b) = - \int_0^{t_0} f'(b e^{\beta(\theta)}) d\theta. \quad (59)$$

Since f' is increasing and $f'(0) = 0$, F_+ is increasing and F_- is decreasing. We also have $F_+(0) = F_-(0) = 0$. Now we fix A, B such that

$$B - A \leq \min\{F_+(a_0), F_-(b_0)\}. \quad (60)$$

We wish to find $a_{A,B} \in [0, a_0], b_{A,B} \in [b_0, 0]$ such that

$$G(a_{A,B}) = G(b_{A,B}) \text{ and } B - A = F_+(a_{A,B}) + F_-(b_{A,B}). \quad (61)$$

Since $B - A \leq \min\{F_+(a_0), F_-(b_0)\}$, by Intermediate Value Theorem there exist $\bar{a} \in [0, a_0], \bar{b} \in [b_0, 0]$ such that $B - A = F_+(\bar{a}) = F_-(\bar{b})$. Define $\lambda := \min\{G(\bar{a}), G(\bar{b})\}$. Without loss of generality, suppose $\lambda = G(\bar{a})$. Then we set $a_1 = \bar{a}$ and $b_1 = 0$. Hence $B - A = F_+(a_1) + F_-(b_1)$. Now we choose $b_2 \in [b_0, 0]$ such that $G(b_2) = G(a_1) \in [0, \lambda]$. Note that $b_2 \geq \bar{b}$ since $G(x)$ is decreasing for $x < 0$ and $G(b_2) \leq \lambda \leq G(\bar{b})$. Since F_- is decreasing, we get $0 \leq F_-(b_2) \leq F_-(\bar{b}) = B - A$. Now by Intermediate Value Theorem, we choose $a_2 \in [0, \bar{a}]$ such that $B - A = F_-(b_2) + F_+(a_2)$. Having defined $\{a_k\}_{1 \leq k \leq n} \subset [0, \bar{a}]$ and $\{b_k\}_{1 \leq k \leq n} \subset [\bar{b}, 0]$ such that $B - A = F_-(b_n) + F_+(a_n)$ we choose $b_{n+1} \in [\bar{b}, 0]$ such that $G(b_{n+1}) = G(a_n)$. Note that the choice of b_{n+1} is guaranteed as $0 \leq G(a_n) \leq \lambda \leq G(\bar{b})$. Subsequently, we get $0 \leq F_-(b_{n+1}) \leq B - A = F_+(\bar{a})$. Now we choose $a_{n+1} \in [0, \bar{a}]$ such that $B - A = F_-(b_{n+1}) + F_+(a_{n+1})$. Hence, by this inductive process we get $\{a_n\}_{n \in \mathbb{N}} \subset [0, \bar{a}]$ and $\{b_n\}_{n \in \mathbb{N}} \subset [\bar{b}, 0]$. Since both sequences are bounded, there is a subsequence n_k such that $b_{n_k} \rightarrow b_{A,B}$ and $a_{n_k} \rightarrow a_{A,B}$ as $k \rightarrow \infty$ for some $a_{A,B} \in [0, \bar{a}]$ and $b_{A,B} \in [\bar{b}, 0]$. Since F_\pm, G are continuous functions, by passing to the limit we show that $a_{A,B} \in [0, \bar{a}]$ and $b_{A,B} \in [\bar{b}, 0]$ satisfy (61).

Suppose $a_{A,B}, b_{A,B}$ are satisfying (61). Now we can choose τ as follows:

$$\tau + \int_0^{t_0} \frac{f(a_{A,B}e^{\beta(\theta)}) - f(b_{A,B}e^{\beta(\theta)})}{a_{A,B} - b_{A,B}} e^{-\beta(\theta)} d\theta = A + \int_0^{t_0} f'(a_{A,B}e^{\beta(\theta)}) d\theta. \quad (62)$$

Since $G(a_{A,B}) = G(b_{A,B})$ and $B - A = F_-(b_{A,B}) + F_+(a_{A,B})$ we get

$$a_{A,B}(\tau - A) + b_{A,B}(B - \tau) = 0. \quad (63)$$

Suppose $u_{A,B}(x, t)$ is the entropy solution to (3) for initial data (57). Then as it has been discussed in [4, section 3], $u_{A,B}$ enjoys the following structure up to time t_0 :

$$u_{A,B}(x, t) = \begin{cases} 0 & \text{if } x \notin [\xi_-(t), \xi_+(t)], \\ \Psi(x - A, t)e^{\beta(t)} & \text{if } \xi_-(t) \leq x \leq \zeta_-(t), \\ a_{A,B}e^{\beta(t)} & \text{if } \zeta_-(t) < x < \zeta_0(t), \\ b_{A,B}e^{\beta(t)} & \text{if } \zeta_0(t) < x < \zeta_+(t), \\ \Psi(x - B, t)e^{\beta(t)} & \text{if } \zeta_+(t) \leq x \leq \xi_+(t), \end{cases} \quad (64)$$

where the curves $\xi_\pm, \zeta_\pm, \zeta_0$ are determined as follows

$$\Psi(\xi_-(t) - A, t) = 0 = \Psi(\xi_+(t) - B, t), \quad (65)$$

$$\Psi(\zeta_-(t) - A, t) = a \text{ and } \Psi(\zeta_+(t) - B, t) = b, \quad (66)$$

$$\zeta_0(t) = \frac{1}{a_{A,B} - b_{A,B}} \int_0^t \left(f(a_{A,B}e^{\beta(\theta)}) - f(b_{A,B}e^{\beta(\theta)}) \right) e^{-\beta(\theta)} d\theta. \quad (67)$$

Note that by hypothesis (H-1) $\xi_-(t) = A$ and $\xi_+(t) = B$. By (62) we have that two curves ζ_\pm meet with each other at point t_0 . For $t \in (t_0, t_0 + \Delta t)$ for small $\Delta t > 0$, the entropy solution $u_{A,B}$ satisfies the following structure

$$u_{A,B}(x, t) = \begin{cases} 0 & \text{if } x \notin [A, B], \\ \Psi(x - A)e^{\beta(t)} & \text{if } A \leq x < \zeta_M(t), \\ \Psi(x - B)e^{\beta(t)} & \text{if } \zeta_M(t) < x \leq B, \end{cases} \quad (68)$$

where $\zeta_M(t)$ is the characteristic curve starting at the point $(\tau, 0)$. Next we claim that $\zeta_M(t) \in (A, B)$ and the structure (68) continues to hold for all $t > t_0$. We can prove this in the same way as it was done for [4, Lemma 3.12]. Suppose the curve $t \mapsto \zeta_M(t)$ intersects either $x = A$ line or $x = B$ line. Without loss of generality we assume that $\zeta_M(t)$ first meets $x = B$ line. Therefore there exists a time $t_1 > 0$ such that at $t = t_1$ we have $\zeta_M(t_1) = B$ and $A < \zeta_M(t) < B$ for $0 \leq t < t_1$. Therefore, (68) is valid up to time t_1 . Consider γ_\pm defined as follows

$$\Psi(\gamma_-(t) - A, t) = \Psi(\zeta_M(t_1) - A, t_1) \text{ and } \gamma_+ = \xi_+. \quad (69)$$

Since (68) is valid up to time t_1 , $\gamma_\pm(t)$ are minimizing curve of the following value function

$$U(x, t) = \min \left\{ \int_0^{\gamma(0)} u_{A,B}^0(y) dy + \int_0^t e^{-\beta(\theta)} f^*(\dot{\gamma}(\theta)) d\theta; \gamma : [0, t] \rightarrow \mathbb{R}, \gamma(t) = x \right\} \quad (70)$$

where f^* is the Legendre transform of f . By (62) we have

$$\int_A^B u_{A,B}^0 = 0. \quad (71)$$

Therefore we have

$$\int_0^{t_1} e^{-\beta(\theta)} f^*(\dot{\gamma}_+(\theta)) d\theta = \int_0^{t_1} e^{-\beta(\theta)} f^*(\dot{\gamma}_-(\theta)) d\theta. \quad (72)$$

By using the definition of Ψ we have

$$\gamma_-(t) - A = \int_0^t f' \left(\Psi(\gamma_-(t) - A, t) e^{\beta(\theta)} \right) d\theta = \int_0^t f' \left(\Psi(\zeta_M(t_1) - A, t_1) e^{\beta(\theta)} \right) d\theta. \quad (73)$$

Differentiating (73) with respect to t , we obtain

$$\dot{\gamma}_-(\theta) = f' \left(\Psi(\zeta_M(t_1) - A, t_1) e^{\beta(\theta)} \right) \text{ for } \theta \in (0, t_1). \quad (74)$$

Similarly, we have

$$\dot{\gamma}_+(\theta) = 0 \text{ for } \theta \in (0, t_1). \quad (75)$$

Since $f^* \geq 0$ we have $f^*(\dot{\gamma}_-(\theta)) = 0$ for a.e. $\theta \in [0, t_1]$. Since 0 is unique minima of f^* , we have $\dot{\gamma}_-(\theta) = 0$ a.e. $\theta \in [0, t_1]$. This gives a contradiction. Hence our claim is proved i.e. $\zeta_M(t) \in (A, B)$ for all $t \geq 0$. Now we are ready to prove Theorem 3.4.

Proof of Theorem 3.4: Define A_n, B_n as follows:

$$A_n = x_n - \frac{1}{n(\log(n+1))^2} \text{ and } B_n = x_n + \frac{1}{n(\log(n+1))^2},$$

$$\text{where } x_n = 4 \sum_{k=1}^{n-1} \frac{1}{k(\log(k+1))^2} + \frac{2}{n(\log(n+1))^2}. \quad (76)$$

From the choice of A_n, B_n and x_n it is clear that

$$B_{n-1} < A_n < B_n < A_{n+1} \text{ and } \lim_{n \rightarrow \infty} x_n = x_0 < \infty. \quad (77)$$

Note that $B_n - A_n = 2n^{-1}(\log(n+1))^{-2} \rightarrow 0$ as $n \rightarrow \infty$. Therefore, there exists $n_0 \in \mathbb{N}$ such that A_n, B_n satisfy (60) for all $n \geq n_0$. By the previous observation, we find $a_{A_n, B_n}, b_{A_n, B_n}$ satisfying (61) for $B_n - A_n$. Next we define initial data u_0 as follows:

$$u_0(x) = \begin{cases} u_{A_n, B_n}^0 & \text{if } x \in [A_n, B_n] \text{ for } n \geq n_0, \\ 0 & \text{otherwise} \end{cases} \quad (78)$$

where u_{A_n, B_n}^0 is defined in (57). To simplify the notation we denote $a_n = a_{A_n, B_n}$ and $b_n = b_{A_n, B_n}$. By using (H-2) in (61) we get

$$B_n - A_n \leq C(a_n - b_n)^q \int_0^{t_0} e^{q\beta(\theta)} d\theta. \quad (79)$$

Therefore, we get

$$a_n - b_n \geq c_0^{-\frac{1}{q}} (B_n - A_n)^{\frac{1}{q}} \text{ where } c_0 = C \int_0^{t_0} e^{q\beta(\theta)} d\theta. \quad (80)$$

From (77) it is clear that u_0 has compact support in \mathbb{R} . By structure (64) and (68) we know that if u_{A_n, B_n} is the entropy solution to (3) for initial data u_{A_n, B_n}^0 then the support of u_{A_n, B_n} lies in the strip $[A_n, B_n] \times [0, \infty)$. Therefore, if $u(x, t)$ is the solution to (3) then u has the following structure

$$u(x, t) = \begin{cases} u_{A_n, B_n}(x, t) & \text{if } x \in [A_n, B_n] \text{ for } n \geq n_0, \\ 0 & \text{otherwise.} \end{cases} \quad (81)$$

By ζ_n we denote the curve ζ_M appeared in the structure (68) of $u_{A, B}$ corresponding to $A = A_n, B = B_n$. From (64) and (68) we obtain the following estimate for any $t > 0$,

$$TV^s(u_{A_n, B_n}(\cdot, t))[A_n, B_n] \geq \min \left\{ (b_n - a_n)^{\frac{1}{s}} e^{\frac{\beta(t)}{s}}, |\Psi(\zeta_n(t) - A_n, t) - \Psi(\zeta_n(t) - B_n, t)|^{\frac{1}{s}} e^{\frac{\beta(t)}{s}} \right\}. \quad (82)$$

From the definition of Ψ we have

$$\zeta_n(t) - A_n - \zeta_n(t) + B_n = \int_0^t \left(f' \left(\Psi(\zeta_n(t) - A_n, t) e^{\beta(\theta)} \right) - f' \left(\Psi(\zeta_n(t) - B_n, t) e^{\beta(\theta)} \right) \right) d\theta. \quad (83)$$

Note that $\zeta_n(t) - A_n > 0 > \zeta_n(t) - B_n$. Since Ψ is increasing in its first variable and $\Psi(0, t) = 0$, we have

$$\Psi(\zeta_n(t) - A_n, t) > 0 > \Psi(\zeta_n(t) - B_n, t). \quad (84)$$

From the decay condition (52) we have

$$B_n - A_n \leq C \int_0^t (\Psi(\zeta_n(t) - A_n, t) - \Psi(\zeta_n(t) - B_n, t))^q e^{q\beta(\theta)} d\theta. \quad (85)$$

Therefore we have

$$|\Psi(\zeta_n(t) - A_n, t) - \Psi(\zeta_n(t) - B_n, t)| \geq (B_n - A_n)^{\frac{1}{q}} \varrho(t)^{-\frac{1}{q}} \text{ where } \varrho(t) := C \int_0^t e^{q\beta(\theta)} d\theta. \quad (86)$$

Combining (80), (82) and (86) we have

$$TV^s(u_{A_n, B_n}(\cdot, t))[A_n, B_n] \geq \min \left\{ c_0^{-\frac{1}{qs}} e^{\frac{1}{s}\beta(t)} (B_n - A_n)^{\frac{1}{qs}}, (B_n - A_n)^{\frac{1}{qs}} \varrho(t)^{-\frac{1}{qs}} e^{\frac{1}{s}\beta(t)} \right\}. \quad (87)$$

Fix an $s > q^{-1}$. Then there exists $\delta > 0$ such that $s = q^{-1} + \delta$. By our choice of A_n and B_n we have

$$B_n - A_n = \frac{2}{n(\log(n+1))^2}. \quad (88)$$

Since $s = (1/q) + \delta$ we have $sq = 1 + q\delta$. We observe that $\beta(t) \geq -t\|\alpha\|_{L^\infty(\mathbb{R}_+)}$ and $\varrho(t) \leq tC e^{q\|\alpha\|_{L^\infty(\mathbb{R}_+)}}$. From definition of c_0 we have $c_0 \leq t_0 C e^{q\|\alpha\|_{L^\infty(\mathbb{R}_+)}}$. Hence, we obtain

$$\begin{aligned} TV^s(u_{A_n, B_n}(\cdot, t))[A_n, B_n] &\geq 2^{\frac{1}{1+q\delta}} e^{\frac{1}{s}\beta(t)} \min \left\{ c_0^{-\frac{1}{qs}}, \varrho(t)^{-\frac{1}{qs}} \right\} n^{-\frac{1}{1+q\delta}} (\log(n+1))^{-\frac{2}{1+q\delta}} \\ &\geq 2^{\frac{1}{1+q\delta}} e^{\frac{-(t+1)}{s}\|\alpha\|_{L^\infty(\mathbb{R}_+)}} \min \left\{ (Ct_0)^{-\frac{1}{qs}}, (Ct)^{-\frac{1}{qs}} \right\} \frac{n^{-\frac{1}{1+q\delta}}}{(\log(n+1))^{\frac{2}{1+q\delta}}}. \end{aligned} \quad (89)$$

Since $[A_n, B_n], n \geq n_0$ are disjoint intervals we have

$$\begin{aligned} TV^s(u(\cdot, t)) &\geq \sum_{n \geq n_0} TV^s(u_{A_n, B_n}(\cdot, t))[A_n, B_n] \\ &\geq 2^{\frac{1}{1+q\delta}} e^{\frac{-(t+1)}{s}\|\alpha\|_{L^\infty(\mathbb{R}_+)}} \min \left\{ (Ct_0)^{-\frac{1}{qs}}, (Ct)^{-\frac{1}{qs}} \right\} \sum_{n \geq n_0} \frac{n^{-\frac{1}{1+q\delta}}}{(\log(n+1))^{\frac{2}{1+q\delta}}} = \infty. \end{aligned}$$

□

4 The scalar multi-D case

In this section we deal with C^∞ -flux function for multi-D scalar conservation laws which reads as follows

$$\partial_t U + \operatorname{div}_X F(U) = 0, \quad U(X, 0) = U_0(X). \quad (90)$$

Non-linearity of a multi-D smooth flux is defined as

Definition 4.1 (Nonlinear flux, [32]). *Let F belong to $C^\infty([a, b], \mathbb{R}^m)$ $[a, b]$ and for each $U \in [a, b]$,*

$$d_F := \sup_{U \in [a, b]} \inf \{k \in \mathbb{N}; k \geq 1, \text{span}(F''(U), \dots, F^{k+1}(U)) = \mathbb{R}^m\} \in \mathbb{N} \cup \{+\infty\}, \quad (91)$$

If $d_F < +\infty$ then F is called a nonlinear flux.

If $d_F = m$ it is called a genuinely nonlinear flux.

It can be checked [32] that since $[a, b]$ is compact $d_F[\cdot]$ attains its maximum at some point $\bar{U} \in [a, b]$, so, d_F is well defined in $\mathbb{N} \cup \{+\infty\}$. By definition of d_F , at least m derivatives of F' are needed to span the m -dimensional space \mathbb{R}^m so $d_F \geq m$. If F is a linear flux, $d_F = +\infty$. Notice that for some exponentially flat fluxes, it is possible to have $d_F = +\infty$ already in dimension one [19, 20]. In this case no BV^s smoothing effect is expected. Indeed, there is a low smoothing effect in a generalized BV space BV^Φ [19, 38, 39].

For a C^∞ nonlinear flux F , the Lions, Perthame and Tadmor conjecture [37] can be reformulated as follows [32], any entropy solution of (90) such that $U_0(\mathbb{R}^m) \subset [a, b]$ are regularized in $W^{s,1}$ for all $s < d_F^{-1}$ where d_F is the non-linearity index as in (91). The Lions, Perthame and Tadmor conjecture is still an open problem.

We prove the limitation of the regularizing effect for the class of C^∞ nonlinear fluxes F such that d_F is odd. The restriction of d_F for odd numbers is due to our previous explicit construction in one dimensional case of solution with the exact maximal regularity for all time only for convex fluxes. The existence of an entropy solution with the conjectured maximal regularity and not more is provided by a construction of a planar wave. This regularity is not improved for large time.

For a bounded strip of time the limitation of the smoothing effect for entropy solutions of multi-dimensional scalar conservation laws in Sobolev spaces has been already proven in [20, 32]. On one hand, the limitation for bounded time was due to the difficulty to study in general behaviour of the solutions after interactions of waves as in [21, 22, 25]. On the other hand, multidimensional fractional BV spaces were not known at that time. Recently, in [28] it has been shown that given a C^2 flux there exists an entropy solution in multi-D such that it is not in BV_{loc} for all time. Authors also prove that there exists an entropy solution which is not in $W_{loc}^{s+\varepsilon,1}, \forall \varepsilon > 0$ for all time with $C^{2,\gamma}$ with $d_F = 1/s$.

The point in this section is to obtain the optimality for all time and in the multi-D BV^s framework. To get the optimality for the multi-D case, a planar wave is used as in [20, 28].

The flux being nonlinear [37] there exist a constant state \underline{U} and a direction ξ such that the flux reaches its degeneracy d_F near \underline{U} and following the direction ξ [32]. That simply means that the scalar flux $f(u) = \xi \cdot F(\underline{U} + u)$ has an exact p -degeneracy (5) with the optimal $p = d_F$. Moreover, for smooth flux, p is an integer [32] bigger than the space dimension: $m \leq p \in \mathbb{N}$. That means that for small u the derivative of the flux f' has exactly a power-law behaviour like u^p . For p odd, f is locally convex (or concave) and the Theorem 3.4 can be used. The result reads as follow.

Corollary 4.1. *Let F be a $C^\infty(\mathbb{R}, \mathbb{R}^m)$ flux with an odd exact degeneracy $d_F = p$ on $[-M, M]$ for some $M > 0$ then there exists an entropy solution U of (90) such that $\forall \varepsilon > 0, \forall t > 0, U(\cdot, t) \in BV_{loc}^s(\mathbb{R}^m, \mathbb{R})$ and $U(\cdot, t) \notin BV_{loc}^{s+\varepsilon}(\mathbb{R}^m, \mathbb{R})$ where $s = 1/p$.*

Remark 4.1. *The parity restriction of p should be neglected with an implicit and more complicated construction used in [28]. Such a solution does not have the same compact support forever.*

We just recall the main features of the proof in [20] for the optimality of the BV^s regularity for a bounded time and then using example of the section 3 and Lemma 2.1 the optimality for all

time follows. First, take an example given in the proof of the Theorem 3.1 with $u_0(x)$ and u the corresponding entropy solution for the flux f and $U_0(X) = \underline{U} + u_0(\xi \cdot X)$ then for all time [20],

$$U(X, t) = \underline{U} + u(\xi \cdot X, t).$$

The BV^s multi-D regularity of the entropy solution U is the consequence of 1-D optimality of u and Lemma 2.1.

5 A class of 2×2 triangular systems

Getting optimal BV^s solutions for general systems for all time is an open problem. Also, the existence of BV^s solutions for systems is in general open. There are some exceptions, BV^s solutions exist for a gas-chromatography system [14], a nonlinear acoustics model and also for diagonal systems [33]. However, the optimality of the regularity is not yet proven. In this section, we consider the first example. The gas-chromatography system is not a Temple system as the well-known chromatography system presented in Bressan's book [15] for instance. Otherwise, this gas-chromatography system enjoys a nice property in Lagrangian variables [14], it has the triangular structure:

$$\partial_t u + \partial_x f(u) = 0 \quad \text{for } x \in \mathbb{R}, t > 0, \quad (92)$$

$$\partial_t v + \partial_x (g(u)v) = 0 \quad \text{for } x \in \mathbb{R}, t > 0, \quad (93)$$

$$(u, v)(x, 0) = (u_0, v_0)(x) \text{ for } x \in \mathbb{R}. \quad (94)$$

At first sight, the system (92)–(93) seems easy to solve. First, one takes the entropy solution of the equation (92). Second, solve the linear transport equation with (93). But, the velocity of the transport equation is $g(u)$ which can be discontinuous. For such equations, a Dirac mass can appear [10]. Thus, due to the transport equation, such systems are not easy to solve in general. The pressureless-gas dynamics system is an example of such problematic systems [11].

In this section we propose optimal BV^s solutions for two cases. First, a self contained construction for a finite time $[0, T_0]$, $T_0 > 0$ where the component v stays continuous. Second, using a recent result of global existence of bounded entropy solutions, we get, as a corollary, the optimality in BV^s for all time.

In the next theorem we construct a solution of the system (92)–(93) such that $(u, v) \notin BV^{s+\varepsilon}(\mathbb{R} \times [0, T_0])$ for all $\varepsilon > 0$ and for power-law type functions f and g satisfying the following relation,

$$g = h \circ f', \quad (95)$$

where h is a Lipschitz function. We first build a continuous solution u to (92) and then solve (93) by using u . Similar line of thought has been previously instrumentalized in [7] to characterize the attainable set for triangular systems.

Theorem 5.1. *Let $T > 0$. Let $f(u) = |u|^{p+1}/(p+1)$, $p \geq 1$, $s = 1/p$ and $g = h \circ f'$ where h is a Lipschitz function. Then there exist compactly supported initial data $u_0, v_0 \in L^\infty(\mathbb{R})$ such that the corresponding entropy solution (u, v) of the triangular system (92)–(93) satisfies $\forall t \in [0, T]$,*

$$u, v \in L^\infty([0, T] \times \mathbb{R}), TV^{s+\varepsilon} u(\cdot, t) = \infty \text{ for all } \varepsilon > 0 \text{ and } TV^{s'} v(\cdot, t) = \infty \text{ for all } s' \in (0, 1].$$

Proof. If $f'(u(x, t))$ is Lipschitz in x -variable then we have $g(u(x, t))$ is Lipschitz in x -variable by the choice of g . The construction is done in two steps

Step 1: Construction of a continuous solution of (1) such that $f'(u(x, t))$ is Lipschitz in the x variable:

Let $\Delta x_n := 1/(n \log^2(n+1))$, $t_n := \frac{\log(n+1)}{\log 2}(T+1) > T$ and let $\delta_n := \left(\frac{\Delta x_n}{t_n}\right)^{\frac{1}{p}}$. We define $x_1 = 0$ and $x_n = 2 \sum_{m=1}^n \Delta x_m$ for $n \geq 2$. Next we consider the following initial data

$$w_0^n(x) = \begin{cases} \left(\frac{\Delta x_n - x}{t_n}\right)^s & \text{if } 0 < x \leq \Delta x_n, \\ -\left(\frac{x - \Delta x_n}{t_n}\right)^s & \text{if } \Delta x_n \leq x < 2\Delta x_n, \\ 0 & \text{otherwise.} \end{cases}$$

Therefore $\forall x \in \mathbb{R}, \forall t > 0$, the entropy solution w_n of (92) with the flux $f(u) = |u|^{p+1}/(p+1)$, is given by

$$w_n(x, t) = \begin{cases} \left(\frac{x}{t}\right)^s & \text{if } 0 < x < \min\{\delta_n^p t, \Delta x_n\}, \\ \left(\frac{\Delta x_n - x}{t_n - t}\right)^s & \text{if } \delta_n^p t < x < \Delta x_n, \\ -\left(\frac{x - \Delta x_n}{t_n - t}\right)^s & \text{if } \Delta x_n < x < 2\Delta x_n - \delta_n^p t, \\ -\left(\frac{2\Delta x_n - x}{t}\right)^s & \text{if } \max\{2\Delta x_n - \delta_n^p t, \Delta x_n\} < x < 2\delta_n, \\ 0 & \text{otherwise.} \end{cases}$$

Let $u_0(x) = \sum_n w_0^n(x - x_n)$ and let u be the entropy solution of (92) with the same flux f . Then one can show that $u(x, t)$ is continuous function on $\mathbb{R} \times (0, T]$ and $f'(u(x, t))$ is Lipschitz in the x variable. We also have

$$TV^{s+\varepsilon} u(\cdot, t)\{[0, 2\Delta x_n]\} \geq 4 (\Delta x_n/t_n)^{\frac{1}{1+p\varepsilon}} \text{ for } t_n > t, \varepsilon > 0.$$

Step 2: We devote this step to find the component v as in (93). In order to do that, it is enough to find a solution of

$$\left. \begin{aligned} \partial_t v + \partial_x(c(x, t)v) &= 0 & \text{for } (x, t) \in \mathbb{R} \times (0, T), \\ v(x, 0) &= v_0(x) & \text{for } x \in \mathbb{R} \end{aligned} \right\} \quad (96)$$

where $c(x, t) = g(u(x, t))$ and u is the entropy solution of (92). We observed that $f'(u)(x, t)$ is Lipschitz in x -variable and so is $c(x, t)$ thanks to (95). We can solve (96) by the method of characteristics and for that, we need to find the solution of the following Cauchy problem

$$\frac{d}{dt} X(t, x_0) = c(X(t, x_0), t), \quad X(0, x_0) = x_0, \quad (97)$$

for each $x_0 \in \mathbb{R}$. By using the classical Cauchy–Lipschitz theorem we obtain a unique solution of (97). In this way, we construct a solution of (96) for L^∞ initial data v_0 . Let v_0 be defined as follows

$$v_0(x) = \begin{cases} -1 & \text{if } 2^{-2k} < x < 2^{-2k+1} \text{ for } k \geq 1, \\ 1 & \text{if } 2^{-2k-1} < x < 2^{-2k} \text{ for } k \geq 1, \\ 1 & \text{if } x > 1/2 \text{ or } x < 0. \end{cases} \quad (98)$$

Consider the sequence $\{y_n\}$ defined as $y_n = (2^{-n} + 2^{-n+1})/2$ for $n \geq 1$. Now fix a $t \in [0, T]$. We define $z_n = X(t, y_n)$. Note that $v(z_n, t) = v_0(y_n)$. Let $s' \in (0, 1]$. By the choice of y_n we get

$$TV^{s'} v(\cdot, t) \geq \sum_{n=1}^{\infty} |v(z_n, t) - v(z_{n+1}, t)|^{1/s'} = \sum_{n=1}^{\infty} |v_0(y_n) - v_0(y_{n+1})|^{1/s'} = \infty.$$

Hence we obtain a solution (u, v) of (92) and (93) such that $u \notin BV^{s+\varepsilon}$ for any $\varepsilon > 0$ and $v \notin BV^{s'}$ for all $s' \in (0, 1]$. \square

Now, the optimality in BV^s for all time is presented. In a recent paper the existence of weak entropy solutions for such triangular system are obtained in [12] under the following assumptions for a convex flux.

- (T-1) The flux $f \in C^4$ is convex and $g \in C^3$.
- (T-2) Initial data u_0 belongs to $BV^{1/3}$ and v_0 to L^∞ .
- (T-3) The system is uniformly strictly hyperbolic,

$$\inf_{|u| \leq M} f' > \sup_{|u| \leq M} g \text{ where } M := \|u_0\|_\infty.$$

Notice, that, in the Theorem 5.1, the system is not assumed to be strictly hyperbolic, which $h = id$ for instance. Here, the strict hyperbolicity is assumed. Moreover, a minimal regularity of the initial data u_0 is needed to ensure the global existence in L^∞ of a solution (u, v) .

Corollary 5.1. *Assume (T-1)–(T-3), and*

$$s = \max \left(\frac{1}{3}, \frac{1}{q} \right),$$

where q is the power of the degeneracy condition (6). There exists an initial data u_0 such that, for all $v_0 \in L^\infty$, the triangular system admits a global solution u staying in BV^s for all time, $v \in L^\infty([0, +\infty), \mathbb{R})$, and, $\forall \epsilon > 0, \forall t > 0, TV^{s+\epsilon} u(\cdot, t) = +\infty$.

6 The multi-D Keyfitz-Kranzer system

In this section, we show that even for data with small total variation, renormalized solution to the Keyfitz-Kranzer system may not be in BV^s . We use the example in [24]. We modified the renormalized solution considered in [24] to show that even if the data has arbitrary small TV the TV^s norm of the solution blows up. Here we mention the key points and the necessary changes. The rest follows from the analysis done in [24].

Consider the following system

$$\begin{aligned}\partial_t u + \operatorname{div}_z(h(|u|)u) &= 0 & \text{for } z \in \mathbb{R}^m, t > 0, \\ u(z, 0) &= u_0(z) & \text{for } z \in \mathbb{R}^m.\end{aligned}\tag{99}$$

where $u : \mathbb{R}^m \times \mathbb{R}_+ \rightarrow \mathbb{R}^k$ and $h \in C^1(\mathbb{R}, \mathbb{R}^m)$. Suppose $\eta := |u|$ solves the following in the sense of Kruřkov

$$\begin{aligned}\partial_t \eta + \operatorname{div}_z(h(\eta)\eta) &= 0 & \text{for } z \in \mathbb{R}^m, t > 0, \\ \eta(\cdot, 0) &= |u_0| & \text{for } z \in \mathbb{R}^m.\end{aligned}\tag{100}$$

Let $\omega := u/|u|$ solves the following transport equation

$$\begin{aligned}\partial_t(\eta\omega) + \operatorname{div}_z(h(\eta)\eta\omega) &= 0 & \text{for } z \in \mathbb{R}^m, t > 0, \\ \omega(\cdot, 0) &= u_0/|u_0| & \text{for } z \in \mathbb{R}^m.\end{aligned}$$

We call $u = \eta\omega$ as *renormalized entropy solution*. Note that the notion of renormalized entropy solution is different from the notion of standard entropy solution. Now we consider a special case of the system (99) with $h = (g, 0, \dots, 0)$. Then we have the following proposition

Proposition 6.1. *Let $h = (g, 0, \dots, 0)$ for $g \in C^1(\mathbb{R})$. Let $k \geq 2, m \geq 2$, and $b \in \mathbb{R}^k \setminus \{0\}$ such that $g'(|b|) \neq 0$. Then there exists a sequence of initial data $u_0^n : \mathbb{R}^m \rightarrow \mathbb{R}^k$ such that*

1. $\|u_0^n - b\|_{BV(\mathbb{R}^m)} + \|u_0^n - b\|_\infty \rightarrow 0$ as $n \rightarrow \infty$,
2. $u_0^n = b$ on $\mathbb{R}^m \setminus B_\lambda(0)$ for some $\lambda > 0$ independent of n ,
3. if u^n is the renormalized entropy solution of (99) with initial data u_0^n then $u^n(\cdot, t) \notin BV_{loc}^s$ for each $n \in \mathbb{N}$, $t \in (0, 1)$ and $s \in (0, 1)$.

Proof. For the sake of simplicity, we prove Proposition 6.1 for $m = 2$ and $k = 2$. Suppose $g'(|b|) = 1$, $g(|b|) = 0$. Let $\delta > 0$ be fixed and $p = s^{-1} > 1$. Let $m_i = i^{p+p\delta}$ for all $i \in \mathbb{N}$. Then we have

$$\sum_i m_i 2^{-i} < +\infty.\tag{101}$$

Let $\varepsilon > 0$ be very small such that

- g is injective on $[|b| - 2\varepsilon, |b| + 2\varepsilon]$,
- $[-\varepsilon, \varepsilon] \subset g([|b| - 2\varepsilon, |b| + 2\varepsilon])$.

Then for sufficiently large i we can choose $r_i \in [-2\varepsilon, 2\varepsilon]$ such that $g(|b| + r_i) = 2^{-i}$. Note that for sufficiently large i we have $r_i \leq 2^{-i+1}$. We write $\beta := b/|b|$, and for each i we choose a $\beta_i \in S^{k-1}$ such that $|\beta - \beta_i| = i^{-1-\delta}$. Consider

$$I_i = [2^{-i}, 2^{-i+1}),\tag{102}$$

$$I_i^j = \left[2^{-i} + \frac{(j-1)2^{-i}}{m_i}, 2^{-i} + \frac{j2^{-i}}{m_i}\right) \quad \text{for } 1 \leq j \leq m_i.\tag{103}$$

Define $\phi_i : \mathbb{R}^2 \rightarrow S^{k-1}$ as follows

$$\phi_i(x, y) := \begin{cases} \beta_i & \text{when } y \in I_i \text{ and } [x2^i] \text{ is odd,} \\ \beta & \text{otherwise.} \end{cases}\tag{104}$$

Also we define $\Lambda_i : \mathbb{R}^2 \rightarrow \mathbb{R}$ as

$$\Lambda_i(x, y) := \begin{cases} r_i & \text{when } y \in I_i \text{ for even } j \text{ and } x \in [-M, M], \\ r_{i+1} & \text{when } y \in I_i \text{ for odd } j \text{ and } x \in [-M, M], \\ 0 & \text{otherwise,} \end{cases} \quad (105)$$

where M is some positive real number bigger than 1 and it will be chosen later. Next we define

$$\eta_0^n := |b| + \sum_{i=n}^{\infty} \Lambda_i, \quad (106)$$

$$\omega_0^n(x, y) := \begin{cases} \phi_i(x, y) & \text{when } y \in I_i \text{ for some } i \geq n \text{ and } x \in [-M, M], \\ \beta & \text{otherwise,} \end{cases} \quad (107)$$

$$u_0^n := \eta_0^n \omega_0^n. \quad (108)$$

Next we show the following two properties of $\{u_0^n\}$ sequence.

M.1 $\|u_0^n - b\|_{L^\infty(\mathbb{R}^2)} \rightarrow 0$ as $n \rightarrow \infty$,

M.2 $\|u_0^n - b\|_{BV(\mathbb{R}^2)} \rightarrow 0$ as $n \rightarrow \infty$.

These have been shown in [24]. For the sake of completion, we briefly mention key steps. From (106)–(108) note that

$$\|u_0^n - b\|_{L^\infty(\mathbb{R}^2)} \leq |b| |\beta - \beta_n| + r_n \leq |b| n^{-1-\delta} + 2^{-n+1} \rightarrow 0 \text{ as } n \rightarrow \infty. \quad (109)$$

We observe that $\text{supp}(u_0^n) \subset [-M, M] \times [0, 1] \subset [-M, M]^2$ as $M > 1$. Note that $\|\eta_0^n\|_{L^\infty(\mathbb{R}^2)} \leq |b| + 1$ and $\|\omega_0^n\|_{L^\infty(\mathbb{R}^2)} \leq 2|\beta| + 1$ for all $n \in \mathbb{N}$. Therefore, to prove (M.2) it is enough to show that

$$\|\eta_0^n - |b|\|_{BV([-2M, 2M]^2)} \rightarrow 0 \text{ and } \|\omega_0^n - \beta\|_{BV([-2M, 2M]^2)} \rightarrow 0 \text{ as } n \rightarrow \infty. \quad (110)$$

Similar to [24], we can show

$$\begin{aligned} \|\eta_0^n - |b|\|_{BV([-2M, 2M]^2)} &\leq 4M^2 \|u_0^n - b\|_{L^\infty(\mathbb{R}^2)} + \sum_{i \geq n} m_i 2^{-i} + (4M + 2)r_n, \\ \|\omega_0^n - \beta\|_{BV([-2M, 2M]^2)} &\leq 4M^2 \|\omega_0^n - \beta\|_{L^\infty(\mathbb{R}^2)} + 2M \sum_{i \geq n} 2^{-i} i^{-1-\delta} 2^i \\ &\quad + 2M \sum_{i \geq n} [i^{-1-\delta} + (i+1)^{-1-\delta}] + (4M + 2)r_n. \end{aligned}$$

Note that $\|\omega_0^n - \beta\|_{L^\infty(\mathbb{R})} \leq |\beta - \beta_n| \rightarrow 0$ as $n \rightarrow \infty$. Since $\sum_{i \geq 1} m_i 2^{-i} < \infty$ and $\sum_{i \geq 1} i^{-1-\delta} < \infty$, we

have

$$\sum_{i \geq n} m_i 2^{-i} \rightarrow 0, \quad \sum_{i \geq n} i^{-1-\delta} \rightarrow 0 \text{ and } \sum_{i \geq n} [i^{-1-\delta} + (i+1)^{-1-\delta}] \rightarrow 0 \text{ as } n \rightarrow \infty.$$

Hence, we obtain (110). Suppose u^n is the unique renormalized solution of (99). We have seen η^n is the unique solution to (100) with initial data η_0^n . Notice that $\eta_0^n(\cdot, y)$ is constant on $[-M, M]$ and by finite speed of conservation laws we get $\eta^n(x, y, t) = \eta_0^n(x, y)$ if $(x, y, t) \in \left\{ \sqrt{x^2 + y^2} \leq C(M - t) \right\}$ where $C = C(\|\eta_0^n\|_\infty)$. Note that for each $R > 0$ we can choose $M > 0$ large enough such that

$$\eta^n(x, y, t) = \eta_0^n(x, y) \text{ for } t \in [0, 1] \text{ and } (x, y) \in (-R, R) \times [0, 1]. \quad (111)$$

We will choose R later. To analyze the angular part $\omega^n := u^n/|u^n|$ we use the fact that η^n is a constant on the curve $\Psi_n(\cdot, x, \cdot)$ where $\Psi_n(\cdot, x, \cdot)$ satisfies

$$\frac{d}{dt}\Psi_n(t, x, y) = h(\eta^n(\Psi_n(t, x, y), t)), \quad (112)$$

$$\Psi_n(0, x, y) = (x, y). \quad (113)$$

We can choose R large enough so that for any $(\tau, x_1, y_1) \in [0, 1]^3 \subset [0, 1] \times [-R, R] \times [0, 1]$, the curve $t \mapsto \Psi_n(t, x_0, y_0)$ lies on the plane $y = y_1$ for $t \in (0, \tau)$ and remains a straight line for $t \in (0, \tau)$ (see [24] for more detailed discussion on this). As it has been observed in [24], choice of R can depend only on g and $\|\rho_0^n\|_{L^\infty(\mathbb{R}^2)}$. Since there exists a constant $C > 0$ such that $\|\rho_0^n\|_{L^\infty(\mathbb{R}^2)} \leq C$ for all $n \geq 1$, we conclude that choice of R does not depend on n . Once we fix the choice of R , we make the choice of M . By a similar discussion as in [24] we have the following,

- if $\eta_0^n(x, y) = |b|$, then $\omega^n(x, y, t) = \omega_0^n(x, y)$,
- if $\eta_0^n(x, y) = |b| + r_i$, then $\omega^n(x, y, t) = \omega_0^n(x - t2^{-i}, y)$.

Therefore, for $j \in \{1, \dots, m_i\}$, $i \geq n$ and $l \in \{1, \dots, 2^i - 1\}$ the function $\omega^n(\cdot, \cdot, t)$ jumps on the segments

$$J_{j,i,l} := \left\{ y = 2^{-i} + \frac{j2^{-i}}{m_i}, x \in [l2^{-i}, (l+t)2^{-i}] \right\}.$$

For a fixed $t > 0$, suppose $\omega^n(\cdot, \cdot, t) \in BV_{loc}^s(\mathbb{R}^2)$. Then, by (11) there exists a $W^n \in BV_{loc}(\mathbb{R}^2)$ and $\pi^n \in Lip^s(\mathbb{R})$ such that $\omega^n(x, y, t) = \pi^n \circ W^n(x, y)$. Hence, $|\omega^n(x_1, y_1, t) - \omega^n(x_2, y_2, t)| \leq C_1 |W^n(x_1, y_1) - W^n(x_2, y_2)|^s$ where $C_1 = Lip^s(\pi^n)$. Let $p = s^{-1}$. Then $|W^n(x_1, y_1) - W^n(x_2, y_2)| \geq C_1^{-p} |\omega^n(x_1, y_1, t) - \omega^n(x_2, y_2, t)|^p$. Now we observe that

$$V_i := \int_{J_{j,i,l}} |W^n(x, y+) - W^n(x, y-)| d\mathcal{H}^1(x) \geq C_1^{-p} t 2^{-i} |\beta - \beta_i|^p = 2^{-i} C_1^{-p} t i^{-p-p\delta},$$

where \mathcal{H}^1 denotes the one dimensional Hausdorff measure. Therefore, we have

$$\|W^n(\cdot, \cdot)\|_{BV([-2M, 2M]^2)} \geq \sum_{i \geq n} \sum_{j=1}^{m_i-1} \sum_{l=1}^{2^i-1} V_i \geq \sum_{i \geq n} \sum_{j=1}^{m_i-1} \sum_{l=1}^{2^i-1} t 2^{-i} i^{-p-p\delta} \geq \frac{t}{2} \sum_{i \geq n} (m_i - 1) i^{-p-p\delta}.$$

Since $m_i = i^{p+p\delta}$ we obtain $\|W^n(\cdot, \cdot)\|_{BV([-2M, 2M]^2)} = \infty$. This gives a contradiction. Hence, $\omega^n(\cdot, \cdot, t) \notin BV_{loc}^s(\mathbb{R}^2)$. \square

Acknowledgement. Authors thank the anonymous referee for valuable comments and suggestions. The first author would like to thank Inspire faculty-research grant DST/INSPIRE/04/2016/000237. Authors would like to thank the IFCAM project ‘‘Conservation laws: BV^s , interface and control’’. The first and third authors acknowledge the support of the Department of Atomic Energy, Government of India, under project no. 12-R&D-TFR-5.01-0520. The second and the fourth authors would like to thank TIFR-CAM for the hospitality.

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